

Regularity of generating families of functions.

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Abstract. We describe the geometric structures involved in the variational formulation of physical theories. In presence of these structures, the constitutive set of a physical system can be generated by a family of functions. We discuss conditions, under which a family of functions generates an immersed Lagrangian submanifold. These conditions are given in terms of the Hessian of the family.

1. Introduction.

The constitutive set of a physical system is frequently a Lagrangian submanifold of a symplectic phase space. Such systems are considered *reciprocal*. It is convenient to be able to derive the constitutive set from a simpler *generating object* such as a Lagrangian in the case of dynamics and an internal energy function in the case of statics. The phase space is not usually the cotangent bundle of a manifold although it is normally isomorphic to a cotangent bundle. We refer to this isomorphism as a *Liouville structure*. For reasons of interpretation the Liouville structure can not be used to replace the phase space by the cotangent bundle. We stress the importance of Liouville structures for variational formulations of physical theories. It is the presence of a Liouville structure that permits the generation of a constitutive set from a generating object. We say that the system is *potential* if its constitutive set is derived from a generating function or a function defined on a *constraint manifold*. Potentiality implies reciprocity. A more general generating object, such as a family of functions does not necessarily generate a Lagrangian submanifold. We discuss sufficient conditions for families of functions to generate Lagrangian submanifolds. We define the *Hessian* of a family of functions at its critical points. The sufficient conditions for families of functions to generate Lagrangian submanifolds are based on this definition.

Reciprocity is an important property of the constitutive set. It can be established by examining directly this set. If the constitutive set is derived from a generating object, then it is more efficient to establish reciprocity by examining the generating object. A similar situation arises when conservation laws are examined. Conservation is a property of dynamics and can be established by direct examination of dynamics. Noether's theorems simplify the procedure by relating conservation properties to invariance properties of the generating object.

The paper is organized as follows. In Section 2 we describe some preliminary constructions. In Sections 3 and 4, we describe geometric structures involved in the variational formulation of physical theories, and the derivation of a set from a generating family of functions. The notion of a critical

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point of a family is introduced. Section 5 contains examples of constitutive sets. In Section 6, we recall results concerning reductions of Lagrangian submanifolds. Then, we discuss the notion of the Hessian of a function (Section 7) and of a family of functions (Section 8), at a critical point. In Section 9 we introduce the notion of a *regular family*, less restrictive then the concept of a Morse family, and we show that the set generated by a regular family is an immersed Lagrangian submanifold.

2. Preliminary constructions.

Let (P, ω) be a symplectic manifold and let V be a vector subspace of the tangent space $T_p P$. We denote by $V^\#$ the *symplectic polar*

$$V^\# = \left\{ \dot{p} \in T_p P; \forall_{\delta p \in V} \langle \omega, \dot{p} \wedge \delta p \rangle = 0 \right\}. \quad (1)$$

If $C \subset P$ is a submanifold, then $T^\# C$ will denote the set

$$\bigcup_{p \in C} (T_p C)^\#. \quad (2)$$

We recall that a submanifold $C \subset P$ is said to be *isotropic* if $T^\# C \supset T C$. A submanifold $C \subset P$ is said to be *coisotropic* if $T^\# C \subset T C$. A submanifold $C \subset P$ is said to be *Lagrangian* if $T^\# C = T C$.

A *symplectic relation* from a symplectic manifold (P_1, ω_1) to a symplectic manifold (P_2, ω_2) is a differential relation ρ from P_1 to P_2 . The graph of a symplectic relation is a Lagrangian submanifold of the symplectic manifold $(P_2 \times P_1, \omega_2 \ominus \omega_1)$. The form $\omega_2 \ominus \omega_1$ is defined by

$$\omega_2 \ominus \omega_1 = pr_2^* \omega_2 - pr_1^* \omega_1, \quad (3)$$

where $pr_1 : P_2 \times P_1 \rightarrow P_1$ and $pr_2 : P_2 \times P_1 \rightarrow P_2$ are the canonical projections.

If C is a coisotropic submanifold of a symplectic manifold (P, ω) , then the set

$$D = \left\{ \dot{p} \in T P; p = \tau_P(\dot{p}) \in C, \forall_{\delta p \in T_p C \subset T_p P} \langle \omega, \dot{p} \wedge \delta p \rangle = 0 \right\} \quad (4)$$

is called the *characteristic distribution* of the symplectic form ω restricted to C . At each $p \in C$ the space $D_p = D \cap T_p P$ is the symplectic polar $T_p^\# C$ of $T_p C$. The characteristic distribution is Frobenius integrable. Its integral manifolds are isotropic submanifolds of (P, ω) called *characteristics* of $\omega|C$. The set of characteristics may be a manifold \tilde{P} . In this case we introduce the *reduction relation* σ from P to \tilde{P} . Its graph is the set

$$\text{graph}(\sigma) = \left\{ (\tilde{p}, p) \in \tilde{P} \times P; p \in \tilde{p} \right\}. \quad (5)$$

Let $\pi : C \rightarrow \tilde{P}$ be the canonical projection. The equality

$$\pi^* \tilde{\omega} = \omega|C \quad (6)$$

defines a symplectic form $\tilde{\omega}$ on \tilde{P} . The reduction relation σ is a symplectic relation from (P, ω) to $(\tilde{P}, \tilde{\omega})$. The graph of σ is the Lagrangian submanifold

$$\text{graph}(\sigma) = \left\{ (\tilde{p}, p) \in \tilde{P} \times P; p \in C, \pi(p) = \tilde{p} \right\}. \quad (7)$$

The projection π is the *strict symplectic reduction* from C onto the symplectic manifold $(\tilde{P}, \tilde{\omega})$ in the terminology of [1]. It is the essential part of the *symplectic reduction relation* σ .

Let F be a function on a differential manifold Q and let $q \in Q$ be a point. The differential of F is a mapping

$$dF: Q \rightarrow T^*Q. \quad (8)$$

At $f = dF(q) \in T^*Q$ we introduce subspaces

$$H_f = TdF(T_q Q) \quad (9)$$

and

$$V_f = \{\delta f \in T_f T^*Q; T\pi_Q(\delta f) = O_{\tau_Q}(q)\} \quad (10)$$

of the vector space $T_f T^*Q$. O_{τ_Q} is the zero section of τ_Q . The intersection $H_f \cap V_f$ of the subspaces is the subspace $\{0\}$ and the sum $H_f + V_f$ is the entire space $T_f T^*Q$. The subspaces H_f and V_f are images of the injections

$$i_h: T_q Q \rightarrow T_f T^*Q: \delta q \mapsto TdF(\delta q) \quad (11)$$

and

$$i_v: T_q^*Q \rightarrow T_f T^*Q: f' \mapsto t Z_{(f, f')}(0), \quad (12)$$

where $Z_{(f, f')}$ is the curve

$$Z_{(f, f')}: \mathbb{R} \rightarrow T^*Q: s \mapsto f + sf'. \quad (13)$$

There are also projections

$$p_h: T_f T^*Q \rightarrow T_q Q: \delta f \mapsto T\pi_Q(\delta f) \quad (14)$$

and

$$p_v: T_f T^*Q \rightarrow T_q^*Q \quad (15)$$

such that the mapping

$$\Psi: T_f T^*Q \rightarrow T_q Q \oplus T_q^*Q: \delta f \mapsto p_h(\delta f) \oplus p_v(\delta f) \quad (16)$$

is the inverse of

$$\Phi: T_q Q \oplus T_q^*Q \rightarrow T_f T^*Q: \delta q \oplus f' \mapsto i_h(\delta q) + i_v(f'). \quad (17)$$

The space $T_f T^*Q$ is a symplectic vector space with a symplectic form

$$\omega_f: T_f T^*Q \times T_f T^*Q \rightarrow \mathbb{R} \quad (18)$$

obtained as a restriction of the symplectic form ω_Q to this vector space. Both subspaces H_f and V_f are Lagrangian subspaces. We choose a pair $(\delta q, f') \in T_q Q \times T_q^*Q$ and use curves

$$\gamma: \mathbb{R} \rightarrow Q \quad (19)$$

and

$$\varphi: \mathbb{R} \rightarrow T^*Q \quad (20)$$

such that $\delta q = t\gamma(0)$, $\varphi(0) = f'$, and $\pi_Q \circ \varphi = \gamma$. The mapping

$$\chi: \mathbb{R}^2 \rightarrow T^*Q: (s_1, s_2) \mapsto Z_{((dF \circ \gamma)(s_1), \varphi(s_1))}(s_2) \quad (21)$$

represents the pair

$$(i_h(\delta q), i_v(f')) \in T_f T^*Q \times T_f T^*Q \quad (22)$$

in the sense that

$$t\chi(\cdot, 0)(0) = i_h(\delta q) \quad (23)$$

and

$$\mathbf{t}\chi(0, \cdot)(0) = i_v(f'). \quad (24)$$

In the following calculation we use the facts that ω_Q is the differential of the Liouville form ϑ_Q , that the Liouville form is vertical and that for each s the curve $\chi(s, \cdot) : \mathbb{R} \rightarrow \mathbf{T}^*Q$ is vertical.

$$\begin{aligned} \omega_f(i_h(\delta q), i_v(f')) &= \langle \omega_Q, i_h(\delta q) \wedge i_v(f') \rangle \\ &= \frac{d}{ds} \langle \vartheta_Q, \mathbf{t}\chi(s, \cdot)(0) \rangle \Big|_{s=0} - \frac{d}{ds} \langle \vartheta_Q, \mathbf{t}\chi(\cdot, s)(0) \rangle \Big|_{s=0} \\ &= -\frac{d}{ds} \langle \vartheta_Q, \mathbf{t}\chi(\cdot, s)(0) \rangle \Big|_{s=0} \\ &= -\frac{d}{ds} \langle \chi(0, s), \mathbf{T}\pi_Q(\mathbf{t}\chi(\cdot, s)(0)) \rangle \Big|_{s=0} \\ &= -\frac{d}{ds} \langle f + sf', \mathbf{t}(\pi_Q \circ \chi(\cdot, s))(0) \rangle \Big|_{s=0} \\ &= -\frac{d}{ds} \langle f + sf', \delta q \rangle \Big|_{s=0} \\ &= -\langle f', \delta q \rangle. \end{aligned} \quad (25)$$

The formula

$$\begin{aligned} \omega_f(i_h(\delta_1 q) + i_v(f'_1), i_h(\delta_2 q) + i_v(f'_2)) &= \omega_f(i_h(\delta_1 q), i_v(f'_2)) + \omega_f(i_v(f'_1), i_h(\delta_2 q)) \\ &= \langle f'_1, \delta_2 q \rangle - \langle f'_2, \delta_1 q \rangle \end{aligned} \quad (26)$$

shows that the mapping (17) is a linear symplectomorphism from the direct product $\mathbf{T}_q Q \oplus \mathbf{T}_q^* Q$ with its canonical symplectic structure to the symplectic vector space $(\mathbf{T}_f \mathbf{T}^* Q, \omega_f)$. The formula

$$\omega_f(\delta_1 f, \delta_2 f) = \langle p_v(\delta_1 f), p_h(\delta_2 f) \rangle - \langle p_v(\delta_2 f), p_h(\delta_1 f) \rangle \quad (27)$$

is equivalent to (26).

3. Subsets of symplectic manifolds generated by families.

The geometric structures involved in the variational formulation of a physical theory are represented by the diagram

$$\begin{array}{ccc} (P, \omega) & \xrightarrow{\alpha} & (\mathbf{T}^* Q, \omega_Q) \\ \pi \downarrow & & \pi_Q \downarrow \\ Q & \xlongequal{\quad} & Q \xleftarrow{\iota} \widetilde{Q} \end{array} \quad \begin{array}{ccc} \overline{Q} & \xrightarrow{\overline{U}} & \mathbb{R} \\ \eta \downarrow & & \\ \widetilde{Q} & & \end{array} \quad (28)$$

The object (P, ω) is the *phase space* of the theory. The diagram

$$\begin{array}{c} (P, \omega) \\ \pi \downarrow \\ Q \end{array} \quad (29)$$

is a vector fibration projecting the phase space onto the *configuration space* Q and the diagram

$$\begin{array}{ccc}
(P, \omega) & \xrightarrow{\alpha} & (\mathsf{T}^*Q, \omega_Q) \\
\pi \downarrow & & \pi_Q \downarrow \\
Q & \xlongequal{\quad} & Q
\end{array} \tag{30}$$

is a vector fibration isomorphism establishing a *Liouville structure* for the phase space (P, ω) . The remaining part

$$\begin{array}{ccc}
\overline{Q} & \xrightarrow{\overline{U}} & \mathbb{R} \\
\eta \downarrow & & \\
Q & \xleftarrow{\iota} & \widetilde{Q}
\end{array} \tag{31}$$

is a *generating object*. It consists of the injection

$$Q \xleftarrow{\iota} \widetilde{Q} \tag{32}$$

of a submanifold $\widetilde{Q} \subset Q$, a differential fibration

$$\begin{array}{c}
\overline{Q} \\
\eta \downarrow \\
\widetilde{Q}
\end{array} \tag{33}$$

and a function $\overline{U} : \overline{Q} \rightarrow \mathbb{R}$ interpreted as a family of functions defined on fibres of the fibration η and denoted by (\overline{U}, η) . The generating object generates a subspace of the phase space. There is an alternate representation of the Liouville structure in terms of a pairing

$$\langle \cdot, \cdot \rangle : P \times_{(\pi, \tau_Q)} \mathsf{T}Q \rightarrow \mathbb{R}. \tag{34}$$

defined by

$$\langle \alpha(p), v \rangle_Q = \langle p, v \rangle \tag{35}$$

for $p \in P$ and each $v \in \mathsf{T}Q$ such that $\tau_Q(v) = \pi(p)$. The canonical pairing

$$\langle \cdot, \cdot \rangle_Q : \mathsf{T}^*Q \times_{(\pi_Q, \tau_Q)} \mathsf{T}Q \rightarrow \mathbb{R} \tag{36}$$

is used. The relation (35) defines the pairing (34) in terms of the symplectomorphism α or the symplectomorphism in terms of the pairing. The set

$$\begin{aligned}
S = & \left\{ p \in P; \tilde{q} = \pi(p) \in \widetilde{Q}, \exists_{\tilde{q} \in \eta^{-1}(\tilde{q})} \forall_{\delta \tilde{q} \in \mathsf{T}_{\tilde{q}} \overline{Q}} \right. \\
& \left. \delta \tilde{q} = \mathsf{T}\eta(\delta \tilde{q}) \Rightarrow \langle p, \delta \tilde{q} \rangle = \langle \mathsf{d}\overline{U}, \delta \tilde{q} \rangle \right\}
\end{aligned} \tag{37}$$

generated by the generating object (31) is expected to be a Lagrangian submanifold of the phase space (P, ω) .

The formula (37) defines the set S directly in terms of the generating object. There is an alternate derivation of this set by the following sequence of operations.

- (1) The function \bar{U} is used to generate the Lagrangian submanifold $\bar{S} = \text{im}(d\bar{U}) \subset T^*\bar{Q}$ of the symplectic manifold $(T^*\bar{Q}, \omega_{\bar{Q}})$.
- (2) The *phase lift* symplectic relation

$$\text{Ph } \eta : (T^*\bar{Q}, \omega_{\bar{Q}}) \rightarrow (T^*\tilde{Q}, \omega_{\tilde{Q}}) \quad (38)$$

of the fibration η is used to produce the set $\tilde{S} = \text{Ph } \eta(\bar{S}) \subset T^*\tilde{Q}$. The relation $\text{Ph } \eta$ can be described in the following way. We denote by $V\bar{Q}$ the subbundle

$$\{\delta\bar{q} \in T\bar{Q}; T\eta(\delta\bar{q}) = 0\} \quad (39)$$

of the tangent bundle $T\bar{Q}$ composed of *vertical* vectors. The *polar*

$$V^\circ\bar{Q} = \left\{ \bar{f} \in T^*\bar{Q}; \forall_{\delta\bar{q} \in V\bar{Q}} \tau_{\bar{Q}}(\delta\bar{q}) = \pi_{\bar{Q}}(\bar{f}) \Rightarrow \langle \bar{f}, \delta\bar{q} \rangle = 0 \right\} \quad (40)$$

of this *vertical subbundle* is a coisotropic submanifold of $(T^*\bar{Q}, \omega_{\bar{Q}})$. Let $\bar{f} \in V^\circ\bar{Q}$, $\bar{q} = \pi_{\bar{Q}}(\bar{f})$, and $q = \eta(\bar{q})$. The relation

$$\langle \tilde{\eta}(\bar{f}), \delta q \rangle = \langle \bar{f}, \delta\bar{q} \rangle \quad (41)$$

with $\delta q \in T_q Q$ and $\delta\bar{q} \in T_{\bar{q}}\bar{Q}$ such that $T\eta(\delta\bar{q}) = \delta q$ defines a differential fibration

$$\tilde{\eta} : V^\circ\bar{Q} \rightarrow T^*\tilde{Q}. \quad (42)$$

This fibration is the strict symplectic reduction (see [1]) from $V^\circ\bar{Q}$ onto the symplectic manifold $(T^*\tilde{Q}, \omega_{\tilde{Q}})$. It is the essential part of the symplectic reduction relation (38) whose graph is the set

$$\left\{ (\tilde{f}, \bar{f}) \in T^*\tilde{Q} \times T^*\bar{Q}; \bar{f} \in V^\circ\bar{Q}, \tilde{f} = \tilde{\eta}(\bar{f}) \right\}. \quad (43)$$

The reduced set

$$\tilde{S} = \text{Ph } \eta(\bar{S}) = \tilde{\eta}(\bar{S} \cap V^\circ\bar{Q}) \quad (44)$$

is not necessarily a Lagrangian submanifold.

- (3) The phase lift

$$\text{Ph } \iota : (T^*\tilde{Q}, \omega_{\tilde{Q}}) \rightarrow (T^*Q, \omega_Q) \quad (45)$$

of the injection ι is applied to the set \tilde{S} . The result is the set $\hat{S} = \text{Ph } \iota(\tilde{S}) \subset T^*Q$. The relation $\text{Ph } \iota$ is, essentially, the strict symplectic reduction from a coisotropic submanifold $\pi_Q^{-1}(\iota(\tilde{Q}))$ of (T^*Q, ω_Q) onto $(T^*\tilde{Q}, \omega_{\tilde{Q}})$. This reduction is the mapping

$$\hat{\iota} : \pi_Q^{-1}(\iota(\tilde{Q})) \rightarrow T^*\tilde{Q} \quad (46)$$

characterized by

$$\langle \hat{\iota}(f), \delta\tilde{q} \rangle = \langle f, T\iota(\delta\tilde{q}) \rangle \quad (47)$$

for each $\delta\tilde{q} \in T_{\tilde{q}}\tilde{Q}$, $\iota(\tilde{q}) = \pi_Q(f)$. If \tilde{S} is a Lagrangian submanifold of $(T^*\tilde{Q}, \omega_{\tilde{Q}})$, then

$$\hat{S} = \text{Ph } \iota(\tilde{S}) = \hat{\iota}^{-1}(\tilde{S} \cap \pi_Q^{-1}(\tilde{Q})) \quad (48)$$

is a Lagrangian submanifold of (T^*Q, ω_Q) ,

(4) The set $S \subset \mathbb{T}^*Q$ is finally obtained as the inverse image $\alpha^{-1}(\widehat{S})$. This set is a Lagrangian submanifold of (P, ω) if \widehat{S} is a Lagrangian submanifold of $(\mathbb{T}^*Q, \omega_Q)$.

In the following example we have a nontrivial Liouville structure and a constrained generating family although the constraint is open.

EXAMPLE 1. Let M be the space time of general relativity with a Minkowski metric $g: \mathbb{T}M \rightarrow \mathbb{T}^*M$ of signature $(1, 3)$. The Lagrangian of a free particle of mass m is the function

$$L: \widetilde{Q} \rightarrow \mathbb{R}: \dot{x} \mapsto m\sqrt{\langle g(\dot{x}), \dot{x} \rangle} \quad (49)$$

defined on the open submanifold

$$\widetilde{Q} = \{\dot{x} \in \mathbb{T}M; \langle g(\dot{x}), \dot{x} \rangle > 0\}$$

of time-like vectors in $Q = \mathbb{T}M$. The dynamics of the particle is a differential equation in the energy-momentum phase space \mathbb{T}^*M . It is therefore a subset $D \subset \mathbb{T}\mathbb{T}^*M$. The space $\mathbb{T}\mathbb{T}^*M$ has a natural symplectic structure. The symplectic form is the total differential $d_T\omega_M$ of the canonical symplectic form ω_M in \mathbb{T}^*M . d_T is a derivation on the exterior algebra of forms on a manifold M with values in the exterior algebra of forms of the tangent bundle $\mathbb{T}M$ (for definition see, e.g., [2]). The dynamics is a Lagrangian submanifold of $(\mathbb{T}\mathbb{T}^*M, d_T\omega_M)$. The Liouville structure

$$\begin{array}{ccc} (\mathbb{T}\mathbb{T}^*M, d_T\omega_M) & \xrightarrow{\alpha_M} & (\mathbb{T}^*\mathbb{T}M, \omega_{\mathbb{T}M}) \\ \mathbb{T}\pi_M \downarrow & & \downarrow \pi_{\mathbb{T}M} \\ \mathbb{T}M & \xlongequal{\quad} & \mathbb{T}M \end{array} \quad (51)$$

is used for generating dynamics from the Lagrangian (49). This Liouville structure was introduced in [3]. It is described rigorously in [4]. At each phase $p \in \mathbb{T}^*M$ the pseudoriemannian structure of M defines subspaces $H_p \subset \mathbb{T}_p \mathbb{T}^*M$ and $V_p \subset \mathbb{T}_p \mathbb{T}^*M$ of horizontal and vertical vectors such that

$$\mathbb{T}_p \mathbb{T}^*M = H_p + V_p \quad (52)$$

$$H_p \cap V_p = \{O_{\mathbb{T}^*M}(p)\}. \quad (53)$$

The dynamics is the set

$$D = \left\{ \dot{p} \in \mathbb{T}\mathbb{T}^*M; \mathbb{T}\pi_M(\dot{p}) \in \widetilde{Q}, \pi_{\mathbb{T}^*M}(\dot{p}) = \frac{mg(\mathbb{T}\pi_M(\dot{p}))}{\|\mathbb{T}\pi_M(\dot{p})\|}, \dot{p} \in H_{\pi_{\mathbb{T}^*M}(\dot{p})} \right\} \quad (54)$$

with

$$\|\mathbb{T}\pi_M(\dot{p})\| = \sqrt{\langle g(\mathbb{T}\pi_M(\dot{p})), \mathbb{T}\pi_M(\dot{p}) \rangle} \quad (55)$$

▲

EXAMPLE 2. The Hamiltonian generating object for the dynamics of Example 1

$$H: \overline{P} \rightarrow \mathbb{R}: (p, \lambda) \mapsto \lambda(\sqrt{\langle p, g^{-1}(p) \rangle} - m) \quad (56)$$

is defined on $\overline{P} = \widetilde{P} \times \mathbb{R}_+$, where \widetilde{P} is the set

$$\widetilde{P} = \{p \in \mathbb{T}^*M; \langle p, g^{-1}(p) \rangle > 0\}$$

is treated as a family of functions on fibres of the projection

$$\zeta : \overline{P} \rightarrow \widetilde{P} : (p, \lambda) \mapsto p. \quad (58)$$

The Liouville structure

$$\begin{array}{ccc} (\mathsf{T}\mathsf{T}^*M, d_T\omega_M) & \xrightarrow{\beta_{(\mathsf{T}\mathsf{T}^*M, \omega_M)}} & (\mathsf{T}^*\mathsf{T}^*M, \omega_{\mathsf{T}^*M}) \\ \tau_{\mathsf{T}^*M} \downarrow & & \pi_{\mathsf{T}^*M} \downarrow \\ \mathsf{T}^*M & \xlongequal{\hspace{1cm}} & \mathsf{T}^*M \end{array} \quad (59)$$

is used. \blacktriangle

EXAMPLE 3. Let M be the space-time manifold of General Relativity. It is a pseudoriemannian manifold of dimension 4 with a metric tensor $g : \mathsf{T}M \rightarrow \mathsf{T}^*M$.

The Lagrangian generating family for the dynamics of a massless particle is the function

$$L : \overline{Q} \rightarrow \mathbb{R} : (\dot{x}, \mu) \mapsto \frac{1}{2\mu} \langle g(\dot{x}), \dot{x} \rangle \quad (60)$$

defined on the space $\overline{Q} = \widetilde{Q} \times \mathbb{R}_+$, where \widetilde{Q} is the tangent bundle $\mathsf{T}M$ with the image of the zero section removed is treated as a family

$$L(\dot{x}, \cdot) : \mathbb{R}_+ \rightarrow \mathbb{R} : \mu \mapsto \frac{1}{2\mu} \langle g(\dot{q}), \dot{q} \rangle \quad (61)$$

of functions on the fibres of the projection

$$\eta : \overline{Q} \rightarrow \widetilde{Q}. \quad (62)$$

The dynamics is the set

$$D = \left\{ \dot{p} \in \mathsf{T}\mathsf{T}^*M; \mathsf{T}\pi_M(\dot{p}) \in \widetilde{Q}, \langle g(\mathsf{T}\pi_M(\dot{p})), \mathsf{T}\pi_M(\dot{p}) \rangle = 0, \right. \\ \left. \exists_{\mu \in \mathbb{R}_+} \tau_{\mathsf{T}^*M}(\dot{p}) = \frac{1}{\mu} g(\mathsf{T}\pi_M(\dot{p})), \dot{p} \in H_{\tau_{\mathsf{T}^*M}(\dot{p})} \right\} \quad (63)$$

\blacktriangle

EXAMPLE 4. The Hamiltonian generating object for the dynamics of Example 3 is the function

$$H : \overline{P} \rightarrow \mathbb{R} : (p, \mu) \mapsto \frac{\mu}{2} \langle p, g^{-1}(p) \rangle \quad (64)$$

defined on $\overline{P} = \widetilde{P} \times \mathbb{R}_+$, where \widetilde{P} is the cotangent bundle T^*M with the image of the zero section removed is treated as a family of functions on fibres of the projection

$$\zeta : \overline{P} \rightarrow \widetilde{P} : (p, \mu) \mapsto p. \quad (65)$$

\blacktriangle

It is obvious that the set S is a Lagrangian submanifold if the first two operations listed above produce a Lagrangian submanifold. For this reason we will concentrate our attention on simpler generating objects with trivial Liouville structures and unconstrained families of functions. Such simple

generating objects are encountered in the theory of partially controlled static systems. Variational formulations of dynamics require the use of nontrivial Liouville structures as is seen in the above example. We will derive conditions sufficient for obtaining Lagrangian submanifolds from the simple generating objects.

4. Families of functions and sets generated by families.

The diagram

$$\begin{array}{ccc} \overline{Q} & \xrightarrow{\overline{U}} & \mathbb{R} \\ \eta \downarrow & & \\ Q & & \end{array} \quad (66)$$

representing a simple generating object is relevant for our analysis. This simple object can be obtained from the diagram (28) by setting $\overline{Q} = Q$ and identifying the symplectic space (P, ω) with (T^*Q, ω_Q) or it can be considered an essential portion of the complete diagram (28).

The set

$$\text{Cr}(\overline{U}, \eta) = \left\{ \overline{q} \in \overline{Q}; \forall_{\delta \overline{q} \in V_{\overline{q}} \overline{Q}} \langle d\overline{U}, \delta \overline{q} \rangle = 0 \right\} \quad (67)$$

is the *critical set* of the family (\overline{U}, η) . Elements of the critical set are *critical points* of (\overline{U}, η) . There is a mapping $\kappa(\overline{U}, \eta) : \text{Cr}(\overline{U}, \eta) \rightarrow T^*Q$ characterized by

$$\langle \kappa(\overline{U}, \eta)(\overline{q}), \delta \tilde{q} \rangle = \langle d\overline{U}, \delta \overline{q} \rangle \quad (68)$$

for each $\delta \tilde{q} \in T_{\eta(\overline{q})} Q$ and each $\delta \overline{q} \in T_{\overline{q}} \overline{Q}$ such that $T\eta(\delta \overline{q}) = \delta \tilde{q}$.

The family of functions (\overline{U}, η) generates a set $S \in T^*Q$. This set is obtained by one of the two following constructions.

- (1) Let $\overline{S} = \text{im}(d\overline{U}) \subset T^*\overline{Q}$ be the Lagrangian submanifold generated by the function \overline{U} . The symplectic relation $\text{Ph} \eta$ applied to \overline{S} produces the set

$$\text{Ph} \eta(\overline{S}) = \tilde{\eta}(\overline{S} \cap V^*\overline{Q}). \quad (69)$$

This is the set S generated by the family (\overline{U}, η) .

- (2) The set S is the image of $\kappa(\overline{U}, \eta)$. The formula

$$S = \left\{ \tilde{f} \in T^*Q; \exists_{\overline{q} \in \overline{Q}} \eta(\overline{q}) = \pi_Q(\tilde{f}) \quad \forall_{\delta \overline{q} \in T_{\overline{q}} \overline{Q}} \langle d\overline{U}, \delta \overline{q} \rangle = \langle \tilde{f}, T\eta(\delta \overline{q}) \rangle \right\} \quad (70)$$

gives an explicit description.

5. Examples.

We give examples of constitutive sets of static systems derived from variational principles applied to families of functions. Variational principles of statics are models for all variational principles of classical physics since at the basis of a variational principle there is a Liouville structure formally identifying the principle with that of a static system. Configuration spaces will be constructed using an affine space Q . The model space is a vector space V of dimension 3 with a Euclidean metric $g : V \rightarrow V^*$.

EXAMPLE 5. A material point with configuration q_2 in the affine space Q is connected to a fixed point q_0 with a rigid rod of length a . A second material point with configuration q_1 is tied elastically to q_2 with a spring of spring constant k . The internal configuration space \overline{Q} is the product $Q \times D$, with

$$D = \{q_2 \in Q; \|q_2 - q_0\| = a\}. \quad (71)$$

The set

$$\mathsf{T}D = \{(q_2, \delta q_2) \in Q \times V; \|q_2 - q_0\| = a, \langle g(q_2 - q_0), \delta q_2 \rangle = 0\} \quad (72)$$

is the tangent bundle of D and the set

$$\mathsf{T}^*D = \{(q_2, f_2) \in D \times V^*; \langle f_2, q_2 - q_0 \rangle = 0\} \quad (73)$$

is chosen to represent the dual of $\mathsf{T}D$. We have the identifications

$$\mathsf{T}\overline{Q} = Q \times V \times \mathsf{T}D \quad (74)$$

and

$$\mathsf{T}^*\overline{Q} = Q \times V^* \times \mathsf{T}^*D. \quad (75)$$

The internal energy

$$\overline{U}: \overline{Q} \rightarrow \mathbb{R}: (q_1, q_2) \mapsto \frac{k}{2} \langle g(q_2 - q_1), q_2 - q_1 \rangle \quad (76)$$

of the system generates the *internal constitutive set*

$$\begin{aligned} \overline{S} = \{(q_1, f_1, q_2, f_2) \in Q \times V^* \times \mathsf{T}^*D; & f_1 = kg(q_1 - q_2), \\ & f_2 - kg(q_2 - q_1) = a^{-2} \langle f_2 - kg(q_2 - q_1), q_2 - q_0 \rangle g(q_2 - q_0)\}. \end{aligned} \quad (77)$$

This set is the image of the differential $d\overline{U}$.

The configuration q_2 is not controlled. The control configuration space is the space Q . The projection

$$\eta: \overline{Q} \rightarrow Q: (q_1, q_2) \mapsto q_1 \quad (78)$$

is the *control relation*. The set

$$\mathsf{V}\overline{Q} = \{(q_1, \delta q_1, q_2, \delta q_2) \in Q \times V \times \mathsf{T}D; \delta q_1 = 0\} \quad (79)$$

is the vertical bundle and the set

$$\mathsf{V}^*\overline{Q} = \{(q_1, f_1, q_2, f_2) \in Q \times V^* \times \mathsf{T}^*D; f_2 = 0\} \quad (80)$$

is its polar. The strict symplectic reduction is the mapping

$$\tilde{\eta}: \mathsf{V}^*\overline{Q} \rightarrow Q \times V^*: (q_1, f_1, q_2, f_2) \mapsto (q_1, f_1). \quad (81)$$

The set

$$\{(q_1, f_1, q_2, f_2) \in Q \times V^* \times \mathsf{T}^*D; f_1 = kg(q_1 - q_2), f_2 = 0, \|q_1 - q_0\|(q_2 - q_0) = \pm(q_1 - q_0)\} \quad (82)$$

is the intersection $\overline{S} \cap \mathsf{V}^*\overline{Q}$. The *constitutive set*

$$\begin{aligned} S = \{(q_1, f_1) \in Q \times V^*; & \|f_1\| = ka \quad \text{if} \quad q_1 = q_0, \\ & f_1 = k(1 \pm a\|q_1 - q_0\|^{-1}) g(q_1 - q_0) \quad \text{if} \quad q_1 \neq q_0\}. \end{aligned} \quad (83)$$

of the partially controlled system is obtained from \overline{S} by applying the symplectic reduction relation $\mathsf{Ph}\eta$. It is the image of $\overline{S} \cap \mathsf{V}^*\overline{Q}$ by the mapping $\tilde{\eta}$.

The internal energy is treated as a family of functions (\overline{U}, η) defined on fibres of the projection η . The critical set

$$\begin{aligned} \mathsf{Cr}(\overline{U}, \eta) = \{(q_1, q_2) \in \overline{Q}; & \|q_1 - q_0\|(q_2 - q_0) = \pm a(q_1 - q_0)\} \\ = \{(q_1, q_2) \in \overline{Q}; & (q_2 - q_1) = a^{-2} \langle g(q_2 - q_1), q_2 - q_0 \rangle (q_2 - q_0)\} \end{aligned} \quad (84)$$

is a submanifold of \overline{Q} . This observation will be confirmed subsequently. The constitutive set S is the image of the injective mapping

$$\kappa(\overline{U}, \eta) : \text{Cr}(\overline{U}, \eta) \rightarrow Q \times V^* : (q_1, q_2) \mapsto (q_1, kg(q_1 - q_2)). \quad (85)$$

The constitutive set can be obtained directly from the variational definition

$$S = \left\{ (q_1, f_1) \in Q \times V^* ; \exists_{q_2 \in D} \forall_{\delta q_1 \in \mathcal{T}Q, \delta q_2 \in \mathcal{T}D} k \langle g(q_2 - q_1), \delta q_2 - \delta q_1 \rangle = \langle f_1, \delta q_1 \rangle \right\}. \quad (86)$$

We show that S is a submanifold of \mathcal{T}^*Q . With the exclusion of the set

$$\{(q_1, f_1) \in Q \times V^* ; q_1 = q_0, \|f_1\| = ka\} \quad (87)$$

the set S is the union of images of the two smooth sections

$$\sigma^+ : Q \setminus \{q_0\} \rightarrow Q \times V^* : q_1 \mapsto (q_1, 1 + a\|q_1 - q_0\|^{-1}g(q_1 - q_0)) \quad (88)$$

and

$$\sigma^- : Q \setminus \{q_0\} \rightarrow Q \times V^* : q_1 \mapsto (q_1, 1 - a\|q_1 - q_0\|^{-1}g(q_1 - q_0)). \quad (89)$$

The set

$$\{(q_1, f_1) \in Q \times V^* ; g(q_1 - q_0) + \|f_1\|^{-1}(a - k^{-1}\|f_1\|)f_1 = 0\} \quad (90)$$

is the set S with the exclusion of

$$\{(q_1, f_1) \in Q \times V^* ; \|q_1 - q_0\| \geq a\}. \quad (91)$$

The set (90) is the image of the smooth section

$$\rho : V^* \rightarrow Q \times V^* : f_1 \mapsto (\|f_1\|^{-1}(k^{-1}\|f_1\| - a)g^{-1}(f_1), f_1) \quad (92)$$

of the canonical projection of $Q \times V^*$ onto V^* . It follows that S is a submanifold of $Q \times V^*$ of dimension 3. \blacktriangle

EXAMPLE 6. A material point with configuration q_1 in the affine space Q is tied elastically to a fixed point q_0 with a spring of spring constant k_1 . A second material point with configuration q_2 is tied elastically to q_1 with a spring of spring constant k_2 and rest length a . The internal configuration space \overline{Q} is the product $Q \times Q$ and the internal energy is the function

$$\overline{U} : Q \times Q \rightarrow \mathbb{R} : (q_1, q_2) \mapsto \frac{k_1}{2} \langle g(q_1 - q_0), q_1 - q_0 \rangle + \frac{k_2}{2} \left(\sqrt{\langle g(q_2 - q_1), q_2 - q_1 \rangle} - a \right)^2. \quad (93)$$

The internal energy generates the internal constitutive set

$$\begin{aligned} \overline{S} = \left\{ (q_1, f_1, q_2, f_2) \in Q \times V^* \times Q \times V^* ; f_1 + f_2 = k_1 g(q_1 - q_0), \right. \\ \left. f_2 = k_2 \left(1 - \frac{a}{\|q_1 - q_2\|} \right) g(q_2 - q_1) \right\}. \end{aligned} \quad (94)$$

The configuration q_2 is not controlled. The control configuration space is the space Q . The projection

$$\eta : \overline{Q} \rightarrow Q : (q_1, q_2) \mapsto q_1 \quad (95)$$

is the *control relation*. The set

$$\mathcal{V}\overline{Q} = \{(q_1, \delta q_1, q_2, \delta q_2) \in Q \times V \times Q \times V ; \delta q_1 = 0\} \quad (96)$$

is the vertical bundle and the set

$$V^\circ \overline{Q} = \{(q_1, f_1, q_2, f_2) \in Q \times V^* \times Q \times V^*; f_2 = 0\} \quad (97)$$

is its polar. The strict symplectic reduction is the mapping

$$\tilde{\eta}: V^\circ \overline{Q} \rightarrow Q \times V^* : (q_1, f_1, q_2, f_2) \mapsto (q_1, f_1). \quad (98)$$

The intersection $\overline{S} \cap V^\circ \overline{Q}$ is the set

$$\{(q_1, f_1, q_2, f_2) \in Q \times V^* \times Q \times V^*; \|q_2 - q_1\| = a, f_1 = k_1 g(q_1 - q_0), f_2 = 0\}. \quad (99)$$

The *constitutive set*

$$S = \{(q_1, f_1) \in Q \times V^*; f_1 = k_1 g(q_1 - q_0)\}. \quad (100)$$

of the partially controlled system is obtained from \overline{S} by applying the symplectic reduction relation $\text{Ph } \eta$. It is the image of $\overline{S} \cap V^\circ \overline{Q}$ by the mapping $\tilde{\eta}$. This constitutive set is the image of the mapping

$$\kappa(\overline{U}, \eta): \text{Cr}(\overline{U}, \eta) \rightarrow Q \times V^* : (q_1, q_2) \mapsto (q_1, k g(q_1 - q_0)), \quad (101)$$

defined on the critical set

$$\text{Cr}(\overline{U}, \eta) = \{(q_1, q_2) \in Q \times Q; \|q_2 - q_1\| = a\}. \quad (102)$$

The constitutive set can also be obtained from the variational construction

$$S = \left\{ (q_1, f_1) \in Q \times V^*; \exists_{q_2 \in D} \forall_{\delta q_1 \in \text{T}Q, \delta q_2 \in \text{T}D} k \langle g(q_2 - q_1), \delta q_2 - \delta q_1 \rangle = \langle f_1, \delta q_1 \rangle \right\}. \quad (103)$$

▲

6. Regular reductions of Lagrangian submanifolds.

Let

$$\tilde{\eta}: \overline{N} \rightarrow \text{T}^* Q \quad (104)$$

be a strict symplectic reduction from a coisotropic submanifold \overline{N} of a symplectic manifold $(\text{T}^* \overline{Q}, \omega_{\overline{Q}})$ onto a symplectic manifold $(\text{T}^* Q, \omega_Q)$ and let \overline{S} be a Lagrangian submanifold of the symplectic manifold $(\text{T}^* \overline{Q}, \omega_{\overline{Q}})$. We are extracting from [1] and [5] the following facts about the reduced set $S = \tilde{\eta}(\overline{S})$. We assume that the intersection of \overline{S} with \overline{N} is not empty.

- (1)) If the intersection of \overline{S} with \overline{N} is clean, then S is an immersed Lagrangian submanifold of $(\text{T}^* Q, \omega_Q)$.
- (2)) If \overline{S} is transverse to \overline{N} , then S is an immersed Lagrangian submanifold of $(\text{T}^* Q, \omega_Q)$ and $\tilde{\eta}|(\overline{N} \cap \overline{S})$ is an immersion.

Recall that submanifolds \overline{S} and \overline{N} have *clean intersection* if $\overline{S} \cap \overline{N} \subset \text{T}^* \overline{Q}$ is a submanifold and

$$\text{T}_{\overline{f}}(\overline{S} \cap \overline{N}) = \text{T}_{\overline{f}}\overline{S} \cap \text{T}_{\overline{f}}\overline{N} \quad (105)$$

at each $\overline{f} \in \overline{S} \cap \overline{N}$. The submanifold \overline{S} is *transverse* to \overline{N} if

$$\text{T}_{\overline{f}}\overline{S} + \text{T}_{\overline{f}}\overline{N} = \text{T}_{\overline{f}}\text{T}^* \overline{Q}. \quad (106)$$

EXAMPLE 7. We use the notation of Example 5. Let $(\mathbf{q}_1, \mathbf{f}_1, \mathbf{q}_2, \mathbf{f}_2) \in \overline{S} \cap V^\circ \overline{Q}$, i.e.

$$f_2 = 0, f_1 = k g(q_1 - q_2), g(q_2 - q_1) = a^{-2} \langle g(q_2 - q_1), q_2 - q_0 \rangle g(q_2 - q_0), \|q_2 - q_0\| = a. \quad (107)$$

We have

$$\mathsf{T}_f \mathsf{T}^* \overline{Q} = \{(\delta q_1, \delta f_1, \delta q_2, \delta f_2) \in V \times V^* \times V \times V^*; \langle g(q_2 - q_0), \delta q_2 \rangle = 0, \langle \delta f_2, q_2 - q_0 \rangle = 0\}, \quad (108)$$

$$\mathsf{T}_f \overline{N} = \{\delta q_1, \delta f_1, \delta q_2, \delta f_2) \in V \times V^* \times V \times V^*; \langle g(q_2 - q_0), \delta q_2 \rangle = 0, \delta f_2 = 0\}, \quad (109)$$

and

$$\begin{aligned} \mathsf{T}_f \overline{S} = & \{(\delta q_1, \delta f_1, \delta q_2, \delta f_2) \in V \times V^* \times V \times V^*; \langle g(q_2 - q_0), \delta q_2 \rangle = 0, \langle \delta f_2, q_2 - q_0 \rangle = 0, \\ & \delta f_1 = kg(\delta q_1 - \delta q_2), \delta f_2 = kg(\delta q_2 - \delta q_1) - a^{-2} \langle kg(\delta q_2 - \delta q_1), q_2 - q_0 \rangle g(q_2 - q_0) \\ & - a^{-2} \langle kg(q_2 - q_1), \delta q_2 \rangle g(q_2 - q_0) - a^{-2} \langle kg(q_2 - q_1), q_2 - q_0 \rangle g(\delta q_2)\} \end{aligned} \quad (110)$$

For every $\delta f = (\delta q_1, \delta f_1, \delta q_2, \delta f_2) \in \mathsf{T}_f \mathsf{T}^* \overline{Q}$, we put

$$\delta_1 f = \left(-\frac{1}{k} g^{-1}(\delta f_2), -\delta f_2, 0, \delta f_2\right) \quad (111)$$

and

$$\delta_2 f = \left(\delta q_1 - \frac{1}{k} g^{-1}(\delta f_2), \delta q_2, \delta f_1 + \delta f_2, 0\right). \quad (112)$$

A direct check shows that $\delta_1 f \in \mathsf{T}_f \overline{S}$, $\delta_2 f \in \mathsf{T}_f \overline{N}$. Since $\delta_1 f + \delta_2 f = \delta f$, we have

$$\mathsf{T}_f \mathsf{T}^* \overline{Q} = \mathsf{T}_f \overline{N} + \mathsf{T}_f \overline{S}. \quad (113)$$

We conclude that \overline{S} is transverse to \overline{N} at f . ▲

EXAMPLE 8. We use the notation of Example 6. Let $f = (q_1, f_1, q_2, f_2) \in \overline{S} \cap \mathsf{V}^\circ \overline{Q}$, i.e. $\|q_1 - q_2\| = a$, $f_1 = k_1 g(q_1 - q_0)$, and $f_2 = 0$. We have $\mathsf{T}_f \mathsf{T}^* \overline{Q} = V \times V^* \times V \times V^*$,

$$\mathsf{T}_f \mathsf{V}^\circ \overline{Q} = \{(\delta q_1, \delta f_1, \delta q_2, \delta f_2) \in V \times V^* \times V \times V^*; f_2 = 0\}, \quad (114)$$

and

$$\begin{aligned} \mathsf{T}_f \overline{S} = & \{(\delta q_1, \delta f_1, \delta q_2, \delta f_2) \in V \times V^* \times V \times V^*; \delta f_1 + \delta f_2 = k_1 g(\delta q_1), \\ & \delta f_2 = -k_2 \frac{a}{\|q_1 - q_2\|^3} \langle g(q_1 - q_2), \delta q_1 - \delta q_2 \rangle g(q_2 - q_1)\} \end{aligned} \quad (115)$$

Since δf_2 is proportional to $g(q_2 - q_1)$, the algebraic sum $\mathsf{T}_f \mathsf{V}^\circ \overline{Q} + \mathsf{T}_f \overline{S}$ is not, for $\dim \mathsf{V} > 1$, equal to $\mathsf{T}_f \mathsf{T}^* \overline{Q}$ and \overline{S} is not transverse to $\mathsf{V}^\circ \overline{Q}$. On the other hand,

$$\overline{S} \cap \mathsf{V}^\circ \overline{Q} = \{(q_1, f_1, q_2, f_2) \in V \times V^* \times V \times V^*; \|q_2 - q_1\| = a, f_1 = k_1 g(q_1 - q_0), f_2 = 0\} \quad (116)$$

is a submanifold, and

$$\begin{aligned} \mathsf{T}_f (\overline{S} \cap \mathsf{V}^\circ \overline{Q}) = & \{(\delta q_1, \delta f_1, \delta q_2, \delta f_2) \in V \times V^* \times V \times V^*; \langle g(q_2 - q_1), \delta q_2 - \delta q_1 \rangle = 0, \\ & \delta f_1 = k_1 g(\delta q_1), \delta f_2 = 0\}. \end{aligned} \quad (117)$$

Comparing (117) with (114) and (115), we establish the equality

$$\mathsf{T}_f (\overline{S} \cap \mathsf{V}^\circ \overline{Q}) = \mathsf{T}_f (\overline{S}) \cap \mathsf{T}_f (\mathsf{V}^\circ \overline{Q}). \quad (118)$$

It follows that \overline{S} and $\mathsf{V}^\circ \overline{Q}$ have clean intersection. ▲

7. The Hessian of a function at a critical point.

Let Q be a differential manifold and let q be a critical point of a function

$$U:Q \rightarrow \mathbb{R}. \quad (119)$$

The image of the differential

$$dU:Q \rightarrow T^*Q \quad (120)$$

is a Lagrangian submanifold $S \subset T^*Q$ of the symplectic space (T^*Q, ω_Q) . It intersects the image of the zero section

$$O_{\pi_Q}:Q \rightarrow T^*Q \quad (121)$$

at $f = dU(q) = O_{\pi_Q}(q)$. The tangent space $S_f = T_f S = TdU(T_q Q)$ is a Lagrangian subspace of the symplectic vector space $(T_f T^*Q, \omega_f)$. We use the decomposition

$$H_f + V_f \quad (122)$$

of the space $T_f T^*Q$ introduced in Section 2. The function $F = 0$ is used. It follows that $dF = O_{\pi_Q}$. The decomposition makes it possible to define a quadratic generating function

$$h:T_q Q \rightarrow \mathbb{R}: \delta q \mapsto \frac{1}{2} \langle p_v(TdU(\delta q)), \delta q \rangle. \quad (123)$$

The *Hessian* of U at the critical point q is the bilinear symmetric function

$$H(U, q): T_q Q \times T_q Q \rightarrow \mathbb{R} \quad (124)$$

defined as the *polarization*

$$\delta h:T_q Q \times T_q Q \rightarrow \mathbb{R}: (\delta_1 q, \delta_2 q) \mapsto h(\delta_1 q + \delta_2 q) - h(\delta_1 q) - h(\delta_2 q) \quad (125)$$

of the quadratic function h . It follows from elementary linear symplectic algebra that the function h is quadratic and its polarization is a symmetric bilinear mapping. The space S_f is generated by h in the sense that

$$S_f = \left\{ \delta f \in T_f T^*Q; \forall_{\delta q \in T_q Q} \langle p_v(\delta f), \delta q \rangle = \delta h(p_h(\delta f), \delta q) \right\}. \quad (126)$$

It follows from this expression for $S_f = TdU(T_q Q)$ that

$$H(U, q)(\delta_1 q, \delta_2 q) = \langle p_v(TdU(\delta_1 q)), \delta_2 q \rangle. \quad (127)$$

A useful expression

$$H(U, q)(\delta_1 q, \delta_2 q) = \omega_f(TdU(\delta_1 q), T O_{\pi_Q}(\delta_2 q)) = \langle \omega_Q, TdU(\delta_1 q) \wedge T O_{\pi_Q}(\delta_2 q) \rangle \quad (128)$$

is derived by using the formula (25).

The image

$$\Psi(S_f) \subset T_q Q \oplus T_q^*Q \quad (129)$$

is a Lagrangian subspace denoted by L_f . This subspace is the graph of the linear mapping

$$\lambda_f:T_q Q \rightarrow T_q^*Q: \delta q \mapsto p_v(TdU(\delta q)) \quad (130)$$

symmetric in the sense that

$$\langle \lambda_f(\delta_1 q), \delta_2 q \rangle = \langle \lambda_f(\delta_2 q), \delta_1 q \rangle. \quad (131)$$

For the Hessian we have the expression

$$\mathsf{H}(U, q)(\delta_1 q, \delta_2 q) = \langle \lambda_f(\delta_1 q), \delta_2 q \rangle. \quad (132)$$

In the following two propositions we are using a critical point q of a function $U: Q \rightarrow \mathbb{R}$, vectors $\delta_1 q$ and $\delta_2 q$ in $\mathsf{T}_q Q$, and a choice of a mapping $\chi: \mathbb{R}^2 \rightarrow Q$ such that $\chi(0, 0) = q$, $\mathsf{t}\chi(\cdot, 0)(0) = \delta_1 q$, and $\mathsf{t}\chi(0, \cdot)(0) = \delta_2 q$.

PROPOSITION 1. *The derivative*

$$\mathsf{D}^{(1,1)}(U \circ \chi)(0, 0) \quad (133)$$

of a function $U: Q \rightarrow \mathbb{R}$ depends on $\delta_1 q$ and $\delta_2 q$ but not on the choice of the mapping χ .

PROOF:

$$\mathsf{D}^{(1,1)}(U \circ \chi)(0, 0) = \mathsf{D}^{(1,1)}((U - U(q)1) \circ \chi)(0, 0) \quad (134)$$

and $U - U(q)1$ is in $\mathsf{I}_1(Q, q) = (\mathsf{I}_0(Q, q))^2$. $\mathsf{I}_0(Q, q)$ is the maximal ideal of functions related to q . It is sufficient to examine the expression (133) for $U = FG$ with F and G in $\mathsf{I}_0(Q, q)$. The equality

$$\begin{aligned} \mathsf{D}^{(1,1)}(FG \circ \chi)(0, 0) &= \mathsf{D}^{(1,1)}((F \circ \chi)(G \circ \chi))(0, 0) \\ &= \mathsf{D}^{(0,1)} \left(\mathsf{D}^{(1,0)}(F \circ \chi) \mathsf{D}^{(0,0)}(G \circ \chi) \right. \\ &\quad \left. + \mathsf{D}^{(0,0)}(F \circ \chi) \mathsf{D}^{(1,0)}(G \circ \chi) \right) (0, 0) \\ &= \left(\mathsf{D}^{(1,1)}(F \circ \chi) \mathsf{D}^{(0,0)}(G \circ \chi) + \mathsf{D}^{(1,0)}(F \circ \chi) \mathsf{D}^{(0,1)}(G \circ \chi) \right. \\ &\quad \left. + \mathsf{D}^{(0,1)}(F \circ \chi) \mathsf{D}^{(1,0)}(G \circ \chi) + \mathsf{D}^{(0,0)}(F \circ \chi) \mathsf{D}^{(1,1)}(G \circ \chi) \right) (0, 0) \\ &= \left(\mathsf{D}^{(1,0)}(F \circ \chi) \mathsf{D}^{(0,1)}(G \circ \chi) + \mathsf{D}^{(0,1)}(F \circ \chi) \mathsf{D}^{(1,0)}(G \circ \chi) \right) (0, 0) \\ &= \langle \mathsf{d}F, \delta_1 q \rangle \langle \mathsf{d}G, \delta_2 q \rangle + \langle \mathsf{d}F, \delta_2 q \rangle \langle \mathsf{d}G, \delta_1 q \rangle \end{aligned} \quad (135)$$

proves the proposition. ■

PROPOSITION 2. *The Hessian $\mathsf{H}(U, q)$ is the bilinear symmetric mapping*

$$(\delta_1 q, \delta_2 q) \mapsto \mathsf{D}^{(1,1)}(U \circ \chi)(0, 0). \quad (136)$$

PROOF: We choose a mapping $\psi: \mathbb{R}^2 \rightarrow \mathsf{T}^* Q$ such that $\pi_Q \circ \psi = \chi$, $\psi(\cdot, 0) = \mathsf{d}U \circ \chi(\cdot, 0)$, and $\psi(0, \cdot) = O_{\pi_Q} \circ \chi(0, \cdot)$. The mapping ψ represents the pair

$$(\mathsf{T}\mathsf{d}U(\delta_1 q), \mathsf{T}O_{\pi_Q}(\delta_2 q)) \in \mathsf{T}_f \mathsf{T}^* Q \times \mathsf{T}_f \mathsf{T}^* Q \quad (137)$$

since

$$\mathsf{T}\mathsf{d}U(\delta_1 q) = \mathsf{T}\mathsf{d}U(\mathsf{t}\chi(\cdot, 0)(0)) = \mathsf{t}(\mathsf{d}U \circ \chi(\cdot, 0))(0) = \mathsf{t}\psi(\cdot, 0)(0) \quad (138)$$

and

$$\mathsf{T}O_{\pi_Q}(\delta_2 q) = \mathsf{T}O_{\pi_Q}(\mathsf{t}\chi(0, \cdot)(0)) = \mathsf{t}(O_{\pi_Q} \circ \chi(0, \cdot))(0) = \mathsf{t}\psi(0, \cdot)(0). \quad (139)$$

The equality

$$\begin{aligned}
\mathsf{H}(U, q)(\delta_1 q, \delta_2 q) &= \langle \omega_Q, \mathsf{T}dU(\delta_1 q) \wedge \mathsf{T}O_{\pi_Q}(\delta_2 q) \rangle \\
&= \frac{d}{ds} \langle \vartheta_Q, \mathsf{t}\psi(s, \cdot)(0) \rangle|_{s=0} - \frac{d}{ds} \langle \vartheta_Q, \mathsf{t}\psi(\cdot, s)(0) \rangle|_{s=0} \\
&= \frac{d}{ds} \langle \tau_{\mathsf{T}^*Q}(\mathsf{t}\psi(s, \cdot)(0)), \mathsf{T}\pi_Q(\mathsf{t}\psi(s, \cdot)(0)) \rangle|_{s=0} \\
&\quad - \frac{d}{ds} \langle \tau_{\mathsf{T}^*Q}(\mathsf{t}\psi(\cdot, s)(0)), \mathsf{T}\pi_Q(\mathsf{t}\psi(\cdot, s)(0)) \rangle|_{s=0} \\
&= \frac{d}{ds} \langle \psi(s, 0), \mathsf{t}(\pi_Q \circ \psi)(s, \cdot)(0) \rangle|_{s=0} - \frac{d}{ds} \langle \psi(0, s), \mathsf{t}(\pi_Q \circ \psi)(\cdot, s)(0) \rangle|_{s=0} \quad (140) \\
&= \frac{d}{ds} \langle (dU \circ \chi)(s, 0), \mathsf{t}\chi(s, \cdot)(0) \rangle|_{s=0} - \frac{d}{ds} \langle (O_\pi \circ \psi)(0, s), \mathsf{t}\chi(\cdot, s)(0) \rangle|_{s=0} \\
&= \frac{d}{ds} \langle (dU(\chi(s, 0)), \mathsf{t}\chi(s, \cdot)(0)) \rangle|_{s=0} \\
&= \frac{\partial}{\partial s} \frac{\partial}{\partial t} U(\chi(s, t))|_{s=0, t=0} \\
&= \mathsf{D}^{(1,1)}(U \circ \chi)(0, 0).
\end{aligned}$$

proves the proposition. \blacksquare

The last proposition offers an alternate definition of the Hessian. This definition is closer to the usual definition of the Hessian in terms of local coordinates.

If q is not a critical point of the function U , then a Hessian of U at q can be defined in relation to a function F on Q such that $dF(q) = dU(q)$. This *relative Hessian* is the Hessian $\mathsf{H}(U - F, q)$.

8. The Hessian of a family of functions at a critical point.

If $\bar{q} \in \bar{Q}$ is a critical point of a family

$$\begin{array}{ccc}
\bar{Q} & \xrightarrow{\bar{U}} & \mathbb{R} \\
\eta \downarrow & & \\
Q & &
\end{array} \quad (141)$$

then $d\bar{U}(\bar{q})$ is in $\mathsf{V}_{\bar{q}}^{\circ}Q$. It follows that

$$\langle d\bar{U}(\bar{q}), \delta\bar{q} \rangle = \langle \tilde{\eta}(d\bar{U}(\bar{q})), \mathsf{T}\eta(\delta\bar{q}) \rangle \quad (142)$$

for each $\delta\bar{q} \in \mathsf{T}_{\bar{q}}\bar{Q}$. Let F be a function on Q such that $dF(\eta(\bar{q})) = \tilde{\eta}(d\bar{U}(\bar{q}))$ and let $\bar{F} = F \circ \eta$. For each $\delta\bar{q} \in \mathsf{T}_{\bar{q}}\bar{Q}$, we have

$$\langle d\bar{F}(\bar{q}), \delta\bar{q} \rangle = \langle (\eta^* dF)(\bar{q}), \delta\bar{q} \rangle = \langle dF(\eta(\bar{q})), \mathsf{T}\eta(\delta\bar{q}) \rangle. \quad (143)$$

Hence, $d\bar{F}(\bar{q}) = d\bar{U}(\bar{q})$. We examine the bilinear mapping

$$\mathsf{V}_{\bar{q}}\bar{Q} \times \mathsf{T}_{\bar{q}}\bar{Q} \rightarrow \mathbb{R} : (\delta_1\bar{q}, \delta_2\bar{q}) \mapsto \mathsf{H}(\bar{U} - \bar{F}, \bar{q})(\delta_1\bar{q}, \delta_2\bar{q}) \quad (144)$$

extracted from the relative Hessian

$$\mathsf{H}(\bar{U} - \bar{F}, \bar{q}) : \mathsf{T}_{\bar{q}}\bar{Q} \times \mathsf{T}_{\bar{q}}\bar{Q} \rightarrow \mathbb{R}. \quad (145)$$

The mapping

$$\chi : \mathbb{R}^2 \rightarrow \bar{Q} \quad (146)$$

representing a pair $(\delta_1 \bar{q}, \delta_2 \bar{q}) \in V_{\bar{q}} \bar{Q} \times T_{\bar{q}} \bar{Q}$ can be chosen to be vertical in the sense that

$$(\eta \circ \chi)(s_1, s_2) = (\eta \circ \chi)(0, s_2). \quad (147)$$

For the function \bar{F} we have

$$(\bar{F} \circ \chi)(s_1, s_2) = (F \circ \eta \circ \chi)(s_1, s_2) = (F \circ \eta \circ \chi)(0, s_2) = (\bar{F} \circ \chi)(0, s_2). \quad (148)$$

It follows that

$$H(\bar{U} - \bar{F}, \bar{q})(\delta_1 \bar{q}, \delta_2 \bar{q}) = D^{(1,1)}((\bar{U} \circ \chi) - (\bar{F} \circ \chi))(0, 0) = D^{(1,1)}(\bar{U} \circ \chi)(0, 0). \quad (149)$$

We had to choose a function F to be able to define the relative Hessian $H(\bar{U} - \bar{F}, \bar{q})$. It turns out that the choice of this function has no effect on the construction of the mapping (144). We define the *Hessian* of the family (141) at the critical point \bar{q} as the bilinear mapping

$$H(\bar{U}, \eta, \bar{q}) : V_{\bar{q}} \bar{Q} \times T_{\bar{q}} \bar{Q} \rightarrow \mathbb{R} : (\delta_1 \bar{q}, \delta_2 \bar{q}) \mapsto H(\bar{U} - \bar{F}, \bar{q})(\delta_1 \bar{q}, \delta_2 \bar{q}). \quad (150)$$

EXAMPLE 9. We consider the generating family of Example 5. Let $\bar{q} = (q_1, q_2) \in Cr(\bar{U}, \eta)$, $\delta_1 \bar{q} = (\delta_1 q_1, \delta_2 q_1) \in T_{\bar{q}} \bar{Q}$, and $\delta_2 \bar{q} = (0, \delta_2 q_2) \in V_{\bar{q}} \bar{Q}$. A mapping $\chi : \mathbb{R}^2 \rightarrow \bar{Q}$ can be chosen of the form

$$\chi(s_1, s_2) = (\chi_1(s_1, s_2), \chi_2(s_1)), \quad (151)$$

where χ_1 represents the pair $\delta_1 q_1, \delta_2 q_1 \in T_{q_1} Q$, and χ_2 represents the vector $\delta_2 q_2 \in T_{q_2} Q$. We have from (149) and (150)

$$H(\bar{U}, \eta, \bar{q})(\delta_1 \bar{q}, \delta_2 \bar{q}) = D^{(1,1)}(\bar{U} \circ \chi)(0, 0) = \langle g(\delta_1 q_2), \delta_2 q_2 - \delta_1 q_1 \rangle. \quad (152)$$

▲

EXAMPLE 10. Here, we consider the generating family of Example 6. At

$$\bar{q} = (q_1, q_2) \in Cr(\bar{U}, \eta) = \{(q_1, q_2) \in Q \times Q; \|q_1 - q_2\| = 0\}, \quad (153)$$

we have

$$H(\bar{U}, \eta, \bar{q})(\delta_1 \bar{q}, \delta_2 \bar{q}) = k_2 \frac{1}{a^2} \langle g(q_2 - q_1), \delta_2 q_2 - \delta_1 q_1 \rangle \langle g(q_2 - q_1), \delta_1 q_2 \rangle. \quad (154)$$

▲

9. Regular families of generating functions.

Let

$$\begin{array}{c} \bar{Q} \\ \downarrow \eta \\ Q \end{array}$$

be a differential fibration, let \bar{q} be a point in \bar{Q} and let \bar{f} be an element of $V_{\bar{q}}^{\circ} \bar{Q}$. We choose a function $F : Q \rightarrow \mathbb{R}$ such that $d_{\eta(\bar{q})} F = \tilde{\eta}(\bar{f})$ and use the function $\bar{F} = F \circ \eta$ to define a splitting $T_{\bar{f}} T^* \bar{Q} = H_{\bar{f}} + V_{\bar{f}}$ at $d\bar{F}(\bar{q}) = \bar{f}$. Note that $d\bar{F}(\bar{Q}) \subset V^{\circ} \bar{Q}$. Hence,

$$i_h(T_{\bar{q}} \bar{Q}) = H_{\bar{f}} = T d\bar{F}(T_{\bar{q}} \bar{Q}) \subset T_{\bar{f}} V^{\circ} \bar{Q}. \quad (155)$$

The equality

$$i_v(V_{\bar{f}}^{\circ}\bar{Q}) = V_{\bar{f}} \cap T_{\bar{f}}V^{\circ}\bar{Q} \quad (156)$$

is a consequence of general properties of the injection i_v . The two equalities (155) and (156) result in

$$\Phi(T_{\bar{q}}\bar{Q} \oplus V_{\bar{q}}^{\circ}\bar{Q}) = T_{\bar{f}}V^{\circ}\bar{Q}. \quad (157)$$

The space $V_{\bar{q}}\bar{Q} \oplus \{0\}$ is the symplectic polar of $T_{\bar{q}}\bar{Q} \oplus V_{\bar{q}}^{\circ}\bar{Q}$ in the symplectic space $T_{\bar{q}}\bar{Q} \oplus T_{\bar{q}}^*\bar{Q}$. Hence,

$$(T_fV^{\circ}\bar{Q})^{\perp} = i_h(V_{\bar{q}}\bar{Q}) = Td\bar{F}(V_{\bar{q}}\bar{Q}). \quad (158)$$

This convenient expression for the symplectic polar is obviously independent of the choice of the function F .

Let $\bar{q} \in Cr(\bar{U}, \eta)$ be a critical point of a family (\bar{U}, η) , let \bar{S} be the Lagrangian submanifold $Td\bar{U}(\bar{Q})$ and let $\bar{f} = d\bar{U}(\bar{q}) \in \bar{S}$. Let F be one of the functions on Q used in Section 8 to define the Hessian $H(\bar{U}, \eta, \bar{q})$ at \bar{q} . The function $\bar{F} = F \circ \eta$ is used to construct an isomorphism

$$\Psi: T_{\bar{f}}T^*Q \rightarrow T_{\bar{q}}\bar{Q} \oplus T_{\bar{q}}^*\bar{Q}. \quad (159)$$

The space $T_f\bar{S} \subset T_{\bar{f}}^*\bar{Q}$ is Lagrangian subspace. Its image

$$L_{\bar{f}} = \Psi(T_f\bar{S}) \subset T_{\bar{q}}\bar{Q} \oplus T_{\bar{q}}^*\bar{Q} \quad (160)$$

is the graph of a symmetric linear mapping

$$\lambda_{\bar{f}}: T_{\bar{q}}\bar{Q} \rightarrow T_{\bar{q}}^*\bar{Q}. \quad (161)$$

We have

$$H(\bar{U} - \bar{F}, \bar{q})(\delta_1\bar{q}, \delta_2\bar{q}) = \langle \lambda_{\bar{f}}(\delta_1\bar{q}), \delta_2\bar{q} \rangle \quad (162)$$

We introduce a rather obvious definition

$$\ker H(\bar{U} - \bar{F}, \bar{q}) = \ker \lambda_{\bar{f}} \quad (163)$$

and a less obvious definition

$$\ker H(\bar{U}, \eta, \bar{q}) = \ker \lambda_{\bar{f}} \cap V_{\bar{q}}\bar{Q}. \quad (164)$$

We have then

$$\ker H(\bar{U}, \eta, \bar{q}) = \{\delta\bar{q} \in V_{\bar{q}}\bar{Q}; i_h(\delta\bar{q}) \in T_f\bar{S}\} \quad (165)$$

and

$$\begin{aligned} i_h(\ker H(\bar{U}, \eta, \bar{q})) &= i_h(V_{\bar{q}}) \cap T_f\bar{S} \\ &= \ker T_f\tilde{\eta} \cap T_f\bar{S} \end{aligned} \quad (166)$$

Consequently,

$$\dim(\ker T_f\tilde{\eta} \cap T_f\bar{S}) = \dim(\ker H(\bar{U}, \eta, \bar{q})) = \dim(V_{\bar{q}}\bar{Q}) - \text{rank } H(\bar{U}, \eta, \bar{q}) \quad (167)$$

DEFINITION 1. A family (\bar{U}, η) is called a *Morse family* if the rank of $H(\bar{U}, \eta, \bar{q})$ is maximal at each $\bar{q} \in Cr(\bar{U}, \eta)$. The family (\bar{U}, η) is said to be *regular* if the critical set $Cr(\bar{U}, \eta)$ is a submanifold of \bar{Q} and the rank of $H(\bar{U}, \eta, \bar{q})$ at each $\bar{q} \in Cr(\bar{U}, \eta)$ is equal to the codimension of $Cr(\bar{U}, \eta)$. \blacktriangle

We will show that a regular family generates a Lagrangian submanifold of T^*Q and that a Morse family is regular.

THEOREM 1. *If (\overline{U}, η) is a regular family, then the image of $\kappa(\overline{U}, \eta)$ is an immersed Lagrangian submanifold of T^*Q .*

PROOF: Let $\bar{q} \in \text{Cr}(\overline{U}, \eta)$ and $\bar{f} = d\overline{U}(\bar{q})$. The rank of $\kappa(\overline{U}, \eta)$ at \bar{q} is equal to

$$\dim(T_{\bar{q}}\text{Cr}(\overline{U}, \eta)) - \dim(\ker(T_{\bar{q}}\kappa(\overline{U}, \eta))). \quad (168)$$

We have

$$\begin{aligned} \ker(T_{\bar{q}}\kappa(\overline{U}, \eta)) &= \ker(T_{\bar{f}}\bar{\eta}) \cap T_{\bar{f}}d\overline{U}(\text{Cr}(\overline{U}, \eta)) \\ &= \ker(T_{\bar{f}}\bar{\eta}) \cap T_{\bar{f}}(\overline{S} \cap V^{\circ}\overline{Q}) \\ &\subset \ker(T_{\bar{f}}\bar{\eta}) \cap (T_{\bar{f}}\overline{S}) \cap T_{\bar{f}}(V^{\circ}\overline{Q}) \\ &= \ker(T_{\bar{f}}\bar{\eta}) \cap T_{\bar{f}}\overline{S}. \end{aligned} \quad (169)$$

It follows from (169) and from (167) that

$$\dim(\ker(T_{\bar{q}}\kappa(\overline{U}, \eta))) \leq \dim(\ker(T_{\bar{f}}\bar{\eta}) \cap T_{\bar{f}}\overline{S}) = \dim V_{\bar{q}}\overline{Q} - \text{rank } H(F, \eta, \bar{q}). \quad (170)$$

Since the family (\overline{U}, η) is regular, $\text{rank } H(\overline{U}, \eta, \bar{q}) = \dim V_{\bar{q}}\overline{Q} + \dim Q - \dim \text{Cr}(\overline{U}, \eta)$ and, consequently,

$$\dim(\ker(T_{\bar{q}}\kappa(\overline{U}, \eta))) \leq \dim \text{Cr}(\overline{U}, \eta) - \dim Q. \quad (171)$$

It follows that

$$\dim(\text{im}(T_{\bar{q}}\kappa(\overline{U}, \eta))) = \dim \text{Cr}(\overline{U}, \eta) - \dim(\ker(T_{\bar{q}}\kappa)) \geq \dim Q. \quad (172)$$

On the other hand, $T_{\bar{q}}\kappa(\overline{U}, \eta)$ is the composition of $T_{\bar{q}}d\overline{U}$, restricted to $T_{\bar{q}}\text{Cr}(\overline{U}, \eta)$, and the strict symplectic reduction $T_{\bar{f}}\bar{\eta}$, which is the essential part of the symplectic reduction relation

$$T_{\bar{f}}\text{Ph}\eta: T_{\bar{f}}T^*\overline{Q} \rightarrow T_{\bar{\eta}(f)}T^*Q. \quad (173)$$

The image $T_{\bar{q}}d\overline{U}(T_{\bar{q}}\text{Cr}(\overline{U}, \eta))$ is an isotropic subspace of $T_{\bar{f}}T^*\overline{Q}$ and, consequently, $\text{im}(T_{\bar{q}}\kappa(\overline{U}, \eta))$ is an isotropic subspace of $T_{\bar{\eta}(f)}T^*Q$. This implies the inequality

$$\dim(\text{im}(T_{\bar{q}}\kappa)) \leq \dim Q, \quad (174)$$

and, consequently,

$$\dim(\text{im}(T_{\bar{q}}\kappa)) = \dim(Q). \quad (175)$$

It follows from the constant rank theorem that $S = \kappa(\overline{U}, \eta)(\text{Cr}(\overline{U}, \rho))$ is an immersed submanifold of T^*Q and $\dim(S) = \dim(Q)$. Since S is isotropic it is Lagrangian. \blacksquare

PROPOSITION 3. *A Morse family is regular.*

PROOF: We have to show that the critical set of a Morse family (\overline{U}, η) is a submanifold of dimension $\dim Q$. Let \bar{q} be a critical point of the family, $\bar{f} = d\overline{U}(\bar{q})$ and

$$\Psi: T_{\bar{f}}T^*\overline{Q} \rightarrow T_{\bar{q}}\overline{Q} \oplus T_{\bar{q}}^*Q \quad (176)$$

the isomorphism constructed with a function on Q as in Section 8. The image $L_{\bar{f}} = \Psi(T_{\bar{f}}\overline{S})$ of $T_{\bar{f}}\overline{S}$ is the graph of a symmetric mapping $\lambda_{\bar{f}}: T_{\bar{q}}\overline{Q} \rightarrow T_{\bar{q}}^*Q$. The rank of the Hessian of (\overline{U}, η) at \bar{q} is the rank of $\lambda_{\bar{f}}$ restricted to $V_{\bar{q}}\overline{Q}$. Let $\lambda_{(\bar{f}, v)}: V_{\bar{q}}\overline{Q} \rightarrow T_{\bar{q}}^*Q$ be this restriction. The dual mapping $\lambda_{(\bar{f}, v)}^*: T_{\bar{q}}\overline{Q} \rightarrow V_{\bar{q}}^*Q$ is of the same rank. Since $\lambda_{\bar{f}}$ is symmetric, $\lambda_{(\bar{f}, v)}^* = \rho_{\bar{q}} \circ \lambda_{\bar{f}}$, where $\rho_{\bar{q}}$ is the restriction of the canonical projection

$$\rho: T^*\overline{Q} \rightarrow V^*\overline{Q} \quad (177)$$

to $T_{\bar{q}}^* \bar{Q}$. The injections i_h, i_v induce injections $i_{(h,\rho)} : T_{\bar{q}} \bar{Q} \rightarrow T_{\rho(f)} V^* \bar{Q}$ and $i_{(v,\rho)} : V_{\bar{q}}^* \bar{Q} \rightarrow T_{\rho(f)} V^* \bar{Q}$ and an isomorphism

$$\Psi_\rho : T_{\rho(f)} V^* \bar{Q} \rightarrow T_{\bar{q}} \bar{Q} \oplus V_{\bar{q}}^* \bar{Q}. \quad (178)$$

With this isomorphism, the mapping $T_{\bar{q}}(\rho \circ d\bar{U})$ is represented by $\rho_{\bar{q}} \circ \lambda_{\bar{f}} = \lambda_{(\bar{f},v)}^*$. The mapping $\rho \circ d\bar{F}$ is the zero section of $V^* \bar{Q}$. It follows that the image of $i_{(h,\rho)}$ is tangent to the zero section. We choose a local trivialization

$$\zeta : V_O^* \bar{Q} \rightarrow V_{\bar{q}}^* \bar{Q} \quad (179)$$

of $V^* \bar{Q}$ in a neighbourhood O of \bar{q} . We have $Cr(\bar{U}, \eta) \cap O = (\zeta \circ \rho \circ d\bar{U})^{-1}(0)$ and $T_{\bar{q}}(\zeta \circ \rho \circ d\bar{U}) : T_{\bar{q}} \bar{Q} \rightarrow V_{\bar{q}}^* \bar{Q}$ coincides with $\lambda_{(\bar{f},v)}^*$. The rank of the Hessian of the family (\bar{U}, η) at \bar{q} is the rank of $\lambda_{(\bar{f},v)}^*$ and consequently, the rank of $T_{\bar{q}}(\zeta \circ \rho \circ d\bar{U})$. It is maximal, hence equal $\dim V_{\bar{q}} \bar{Q} = \dim V_{\bar{q}}^* \bar{Q}$. It follows that $T_{\bar{q}}(\zeta \circ \rho \circ d\bar{U})$ is surjective and, by the implicit function theorem, $Cr(\bar{U}, \eta)$ is a submanifold of dimension $\dim \bar{Q} - \dim V_{\bar{q}}^* \bar{Q} = \dim Q$. ■

PROPOSITION 4. *The family (\bar{U}, η) is regular if and only if $\bar{S} = d\bar{U}(\bar{Q})$ and $V^* \bar{Q}$ have clean intersection.*

PROOF: Let $\bar{f} \in \bar{S} \cap V^* \bar{Q}$ and $\pi_{\bar{Q}}(\bar{f}) = \bar{q}$. As in the preceding proposition, we shall use the canonical projection (177) and the isomorphism (178). We have

$$\Psi_\rho(T_{\bar{f}}(T_{\bar{f}} V^* \bar{Q})) = T_{\bar{q}} \bar{Q} \oplus 0 \quad (180)$$

and $\Psi_\rho(T_{\bar{q}}(\rho \circ d\bar{U}))$ is the graph of $\rho_{\bar{q}} \circ \lambda_{\bar{f}} = \lambda_{(\bar{f},v)}^* : T_{\bar{q}} \bar{Q} \rightarrow V_{\bar{q}}^*$. The rank of $\lambda_{(\bar{f},v)}$ is equal to the rank of the Hessian of the family (\bar{Q}, η) at \bar{q} . It follows that

$$T_{\bar{f}}(T_{\bar{f}} V^* \bar{Q} + T_{\bar{f}} \bar{S}) = T_{\bar{q}} \bar{Q} \oplus \text{im}(\lambda_{(\bar{f},v)}), \quad (181)$$

and the dimension of these spaces is $\dim \bar{Q} + \dim \text{im}(\lambda_{\bar{f}}) = \dim \bar{Q} + \text{rank } H(\bar{U}, \eta, \bar{q})$. Since the kernel of $T_{\bar{f}} \rho$ is contained in $T_{\bar{f}} V^* \bar{Q}$, we have

$$\begin{aligned} \dim(T_{\bar{f}} V^* \bar{Q} + T_{\bar{f}} \bar{S}) &= \dim(T_{\bar{f}}(T_{\bar{f}} V^* \bar{Q} + T_{\bar{f}} \bar{S})) + \dim \ker(T_{\bar{f}} \rho) \\ &= \dim \bar{Q} + \text{rank } H(\bar{U}, \eta, \bar{q}) + \dim Q \end{aligned} \quad (182)$$

It follows that

$$\begin{aligned} \dim(T_{\bar{f}} V^* \bar{Q} \cap T_{\bar{f}} \bar{S}) &= \dim(T_{\bar{f}} V^* \bar{Q}) + \dim(T_{\bar{f}} \bar{S}) - \dim(T_{\bar{f}} V^* \bar{Q} + T_{\bar{f}} \bar{S}) \\ &= \dim(\bar{Q}) + \dim(Q) + \dim(\bar{Q}) - \dim(\bar{Q}) - \dim(Q) - \text{rank } H(\bar{U}, \eta, \bar{q}) \\ &= \dim(\bar{Q}) - \text{rank } H(\bar{U}, \eta, \bar{q}) \end{aligned} \quad (183)$$

We conclude that $T_{\bar{f}} Cr(\bar{U}, \eta) = T_{\bar{f}} V^* \bar{Q} \cap T_{\bar{f}} \bar{S}$ if and only if $\dim(Cr(\bar{U}, \eta)) = \dim(\bar{Q}) - \text{rank } H(\bar{U}, \eta, \bar{q})$. ■

COROLLARY 1. *(\bar{U}, η) is a Morse family if and only if $V^* \bar{Q}$ and \bar{S} have transversal intersection.*

PROOF: We have from (182) that $\dim(T_{\bar{f}} V^* \bar{Q} + T_{\bar{f}} \bar{S}) = \dim T^* \bar{Q}$ if and only if

$$\dim(Q) + \text{rank } H(\bar{U}, \eta, \bar{q}) = \dim(\bar{Q}), \quad (184)$$

i.e., if and only if $H(\bar{U}, \eta, \bar{q})$ is of maximal rank. ■

EXAMPLE 11. Let (\bar{U}, η) be the generating family of Example 5. The Hessian of this family

$$H(\bar{U}, \eta, \bar{q})(\delta_1 \bar{q}, \delta_2 \bar{q}) = D^{(1,1)}(\bar{U} \circ \chi)(0, 0) = \langle g(\delta_1 q_2), \delta_2 q_2 - \delta_2 q_1 \rangle \quad (185)$$

is of maximal rank. The family is a Morse family. \blacktriangle

EXAMPLE 12. For the family (\bar{U}, η) of Example 6, the critical set

$$Cr(\bar{U}, \eta) = \{(q_1, q_2) \in Q \times Q; \|q_2 - q_1\| = a\} \quad (186)$$

is a submanifold of codimension 1. The Hessian

$$H(\bar{U}, \eta, \bar{q})(\delta_1 \bar{q}, \delta_2 \bar{q}) = k_2 \frac{1}{a^2} \langle g(q_2 - q_1), \delta_2 q_2 - \delta_2 q_1 \rangle \langle g(q_2 - q_1), \delta_1 q_2 \rangle \quad (187)$$

is of constant rank 1. The family is regular. \blacktriangle

EXAMPLE 13. Let $\bar{Q} = \mathbb{R}^2$, $Q = \mathbb{R}$ and $\eta: \mathbb{R}^2 \rightarrow \mathbb{R}: (x, \lambda) \mapsto x$. For $\bar{U}(x, \lambda) = \lambda x^2$ we have $Cr(\bar{U}, \eta) = \{(x, \lambda): x = 0\}$ and the Hessian is the trivial zero form. In this case the intersection of \bar{S} and $V^{\circ} \bar{Q}$ is not clean, but the Hessian is of constant rank. The generated set is an isotropic submanifold, but not Lagrangian. \blacktriangle

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